ANSWER on Question #81519 - Math - Differential Equations

QUESTION

Reduce the following PDE to a set of three ODEs by the method of separation of variables

$$\frac{1}{r}\frac{\partial}{\partial r}\left(r\cdot\left(\frac{\partial V}{\partial r}\right)\right) + \left(\frac{1}{r^2}\right)\cdot\frac{\partial^2 V}{\partial \theta^2} + \frac{\partial^2 V}{\partial z^2} + k^2\cdot V = 0$$

SOLUTION

Assume

$$V(r, \theta, z) = R(r)\Theta(\theta)Z(z)$$

Then,

$$\begin{aligned}
\Theta Z \frac{1}{r} \frac{\partial}{\partial r} \left(r \cdot \left(\frac{\partial R}{\partial r} \right) \right) + R Z \left(\frac{1}{r^2} \right) \cdot \frac{\partial^2 \Theta}{\partial \theta^2} + R \Theta \frac{\partial^2 Z}{\partial z^2} + k^2 \cdot R \Theta Z = 0 \quad | \quad \div \left(\frac{1}{R \Theta Z} \right) \rightarrow \\
\frac{1}{R} \frac{1}{r} \frac{\partial}{\partial r} \left(r \cdot \left(\frac{\partial R}{\partial r} \right) \right) + \frac{1}{\Theta} \left(\frac{1}{r^2} \right) \cdot \frac{\partial^2 \Theta}{\partial \theta^2} + \frac{1}{Z} \frac{\partial^2 Z}{\partial z^2} + k^2 = 0 \rightarrow \\
\frac{1}{R} \frac{1}{r} \frac{\partial}{\partial r} \left(r \cdot \left(\frac{\partial R}{\partial r} \right) \right) + \frac{1}{\Theta} \left(\frac{1}{r^2} \right) \cdot \frac{\partial^2 \Theta}{\partial \theta^2} + k^2 = -\frac{1}{Z} \frac{\partial^2 Z}{\partial z^2}
\end{aligned}$$

In the above equation the left-hand side depends on r and θ , while the right-hand side depends on z. The only way these two members are going to be equal for all values of r, θ and z is when both of them are equal to a constant. Let us define such constant as $-l^2$.

With this choice for the constant, we obtain:

$$\frac{d^2Z}{dz^2} - l^2 \cdot Z = 0$$

The general solution of this equation is:

$$Z(z) = A_1 e^{lz} + A_2 e^{-lz}$$

Such a solution, when considering the specific boundary conditions, will allow Z(z) to go to zero for z going to $\pm \infty$, which makes physical sense. If we had given the constant a value of l^2 , we would have had periodic trigonometric functions, which do not tend to zero for z going to infinity.

Once sorted the z —dependency, we need take care of r and θ .

$$\frac{1}{R} \frac{1}{r} \frac{\partial}{\partial r} \left(r \cdot \left(\frac{\partial R}{\partial r} \right) \right) + \frac{1}{\Theta} \left(\frac{1}{r^2} \right) \cdot \frac{\partial^2 \Theta}{\partial \theta^2} + k^2 = -l^2 \rightarrow$$

$$\frac{1}{R} \frac{1}{r} \frac{\partial}{\partial r} \left(r \cdot \left(\frac{\partial R}{\partial r} \right) \right) + \frac{1}{\Theta} \left(\frac{1}{r^2} \right) \cdot \frac{\partial^2 \Theta}{\partial \theta^2} = -(k^2 + l^2) \rightarrow$$

$$\frac{1}{R} \frac{1}{r} \left(\frac{\partial R}{\partial r} + r \cdot \frac{\partial^2 R}{\partial r^2} \right) + \frac{1}{\Theta} \left(\frac{1}{r^2} \right) \cdot \frac{\partial^2 \Theta}{\partial \theta^2} = -(k^2 + l^2) \rightarrow$$

$$\frac{1}{R} \frac{\partial^2 R}{\partial r^2} + \frac{1}{R} \frac{\partial R}{\partial r} + \frac{1}{\Theta} \left(\frac{1}{r^2} \right) \cdot \frac{\partial^2 \Theta}{\partial \theta^2} = -(k^2 + l^2) \times (r^2) \rightarrow$$

$$\frac{r^2}{R} \frac{\partial^2 R}{\partial r^2} + \frac{r}{R} \frac{\partial R}{\partial r} + (k^2 + l^2)r^2 = \frac{1}{\Theta} \cdot \frac{\partial^2 \Theta}{\partial \theta^2}$$

Again we are in a situation where the only way a solution can be found for the above equation is when both members are equal to a constant. This time we select a positive constant, which we call m^2 . The equation for Θ becomes, for:

$$\frac{d^2\Theta}{d\theta^2} + m^2\Theta = 0$$

This solution is well suited, to describe the variation for an angular coordinate like θ . Had we chosen the set both members equal to a negative number, we would have ended up with exponential functions with a different value assigned to $\Theta(\theta)$ for each 360° turn, a clear non-physical solution.

Last to be examined is the r –dependency. We have:

$$\frac{r^2}{R}\frac{d^2R}{dr^2} + \frac{r}{R}\frac{dR}{dr} + (k^2 + l^2)r^2 = m^2 \to$$

$$r^2 \cdot \frac{d^2R}{dr^2} + r \cdot \frac{dR}{dr} + [(k^2 + l^2)r^2 - m^2] \cdot R = 0$$

This equation is a well-known equation of mathematical physics called parametric Bessel's equation. With sample linear transformation of variable, $x = r \cdot \sqrt{(k^2 + l^2)}$, equation is readily changed into a Bessel's equation:

$$\frac{dR}{dr} = \frac{dR}{dx} \cdot \frac{dx}{dr} = \sqrt{k^2 + l^2} \cdot R'$$

$$\frac{d^2R}{dr^2} = \frac{d}{dx} \left(\sqrt{k^2 + l^2} \cdot R' \right) \cdot \frac{dx}{dr} = (k^2 + l^2) \cdot R''$$

Then,

$$r^{2} \cdot \frac{d^{2}R}{dr^{2}} + r \cdot \frac{dR}{dr} + \left[(k^{2} + l^{2})r^{2} - m^{2} \right] \cdot R = 0 \to$$

$$\frac{x^{2}}{(k^{2} + l^{2})} \cdot (k^{2} + l^{2}) \cdot R'' + \frac{x}{\sqrt{k^{2} + l^{2}}} \cdot \sqrt{k^{2} + l^{2}} \cdot R' + \left[x^{2} - m^{2} \right] \cdot R = 0 \to$$

$$x^{2} \cdot R'' + x \cdot R' + \left[x^{2} - m^{2} \right] \cdot R = 0$$

where R'' and R' indicate the first and second derivatives with respect to x.

In what follows we will assume that m is a real, non-negative number.

Linearly independent solutions are typically denote $J_m(x)$ (Bessel Functions) and $N_m(x)$ (Neumann Functions).

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